

Relativistic Dynamics:

In relativistic mechanics, the dynamical quantities are: mass, energy, momentum and force, from the point of view of special theory of relativity.

In Newtonian mechanics we have:

$$\vec{P} = m \vec{v}, K = \frac{1}{2} m v^2, F = ma, F = \frac{d\vec{P}}{dt}$$

Also $m = \text{constant}$ that does not change for moving bodies.

What are the corresponding relations in the theory of relativity?

Newtonian Mechanics

- * $\vec{P} = m \vec{v} = m \frac{d\vec{r}}{dt}, K = \frac{1}{2} m v^2 = \frac{\vec{P}^2}{2m}$

- * Vectors have 3 components

- * $\vec{r} = x \hat{i} + y \hat{j} + z \hat{k}$ } Cartesian coordinates
 $\vec{v} = v_x \hat{i} + v_y \hat{j} + v_z \hat{k}$ } (x, y, z)

- * mass : Absolute (quantity of matter).

Independent of velocity.

- * Time : Absolute in different inertial frames

Einstein's Mechanics

- * Vectors have 4 components (x, y, z, t) .

- * mass : Not absolute (depends on velocity).

- * Time : Not absolute (depends on velocity).

Example: $x' = \gamma(x - vt)$

$$y' = y$$

$$z' = z$$

$$t' = \gamma(t - \frac{x_0}{c^2})$$

Now we know that $m = \frac{m_0}{\sqrt{1 - \frac{v^2}{c^2}}} = \gamma m_0$.

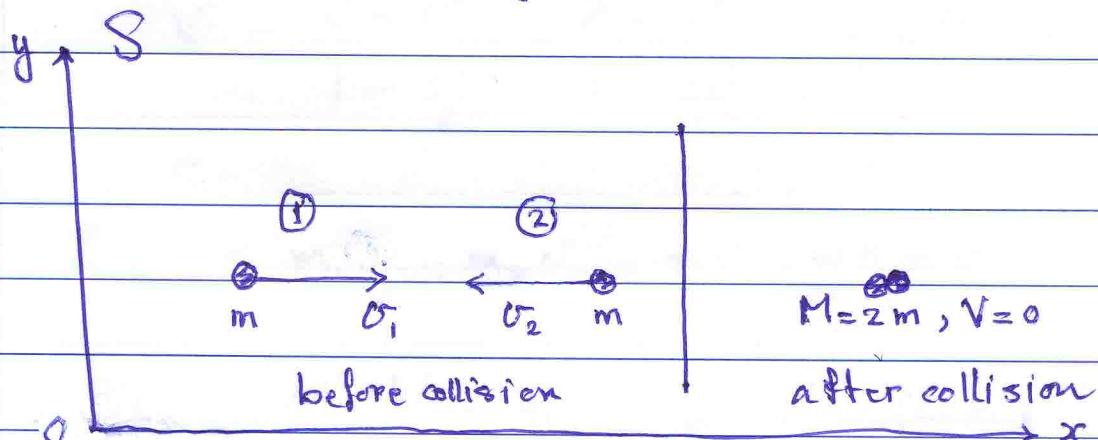
m_0 = rest mass.

m = mass of the object moving with velocity v .

Prove the above relation through an example

Consider a completely inelastic collision between two masses (equal) moving along x -direction with velocities v and $-v$.

In the stationary frame S



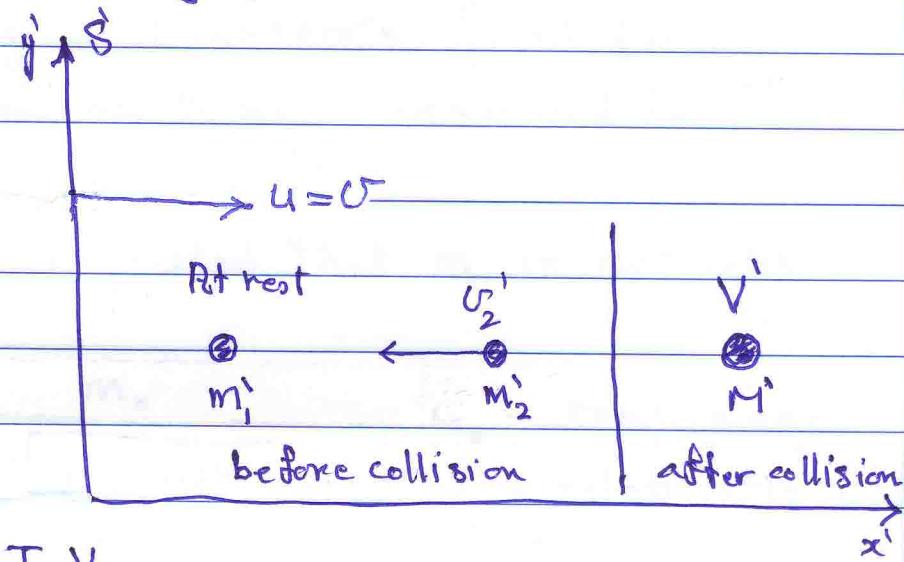
$$|v_1| = |v_2| = v$$

$$v_2 = -v_1 = -v$$

Linear momentum is conserved

$$m v_1 + m v_2 = 2m(0)$$

In the moving frame S'



Using L.T.V.

$$U'_1 = \frac{U_1 - U}{1 - \frac{U_1 U}{c^2}} = \frac{U - U}{1 - \frac{U^2}{c^2}} = 0 \quad (U_1 = U)$$

$$\begin{aligned} U'_2 &= \frac{U_2 - U}{1 - \frac{U_2 U}{c^2}} = \frac{-U - U}{1 + \frac{U^2}{c^2}} \quad (U_2 = -U) \\ &= \frac{-2U}{1 + U^2/c^2} \end{aligned}$$

$$V = \frac{V - U}{1 - \frac{VU}{c^2}} = \frac{0 - U}{1 - \frac{(0)U}{c^2}} = -U \quad (V = 0)$$

\Rightarrow In S :

$$P_{\text{before}} = m_1 U'_1 + m_2 U'_2 = m U - m U = 0$$

$$P_{\text{after}} = (2m)V = 0 \quad \checkmark \quad (\text{OK})$$

In S' :

$$P'_{\text{before}} = m_1 U'_1 + m_2 U'_2 = m(0) + m\left(\frac{-2U}{1 + U^2/c^2}\right)$$

$$P'_{\text{after}} = 2mV = -2mU$$

$$\Rightarrow P'_{\text{before}} \neq P'_{\text{after}} \quad !!$$

According to Einstein's first postulate, linear momentum must be conserved in S'

Einstein suggested that m is not absolute \Rightarrow

$$m = \frac{m_0}{\sqrt{1 - \frac{v^2}{c^2}}}$$

m_0 = rest mass

m = relativistic mass

If this suggestion leads to conservation of linear momentum in S' , then it is acceptable.

We have, in our problem

$$m_1 = \frac{m_0}{\sqrt{1 - \frac{v^2}{c^2}}}, \quad m_2 = \frac{m_0}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$U_1 = U_2 = v$$

$$M = m_1 + m_2 = \frac{2m_0}{\sqrt{1 - \frac{v^2}{c^2}}}$$

Since M is at rest in $S \Rightarrow M = M_0 \quad \text{---} \star$

In S' , m_1 is at rest $\Rightarrow m'_1 = m_0$

m'_2 is moving with $U'_2 = \frac{-2v}{(1 + \frac{v^2}{c^2})}$

$$m'_2 = \frac{m_0}{\sqrt{1 - \frac{1}{c^2} \left(\frac{-2v}{1 + \frac{v^2}{c^2}} \right)^2}}$$

$$= \frac{m_0}{\sqrt{1 - \frac{4v^2}{c^2(1 + \frac{v^2}{c^2})^2}}}$$

$$= \frac{m_0}{\sqrt{1 - \frac{4v^2 c^2}{(c^2 + v^2)^2}}}$$

$$= \frac{m_0}{\sqrt{\frac{(c^2 + v^2)^2 - 4v^2 c^2}{(c^2 + v^2)^2}}}$$

$$= \frac{m_0}{\sqrt{\left(\frac{c^2 - v^2}{c^2 + v^2}\right)^2}}$$

$$= m_0 \left(\frac{1 + \frac{v^2}{c^2}}{1 - \frac{v^2}{c^2}} \right)$$

The combined mass M' is moving with $V' = -v$,
so

$$M' = \frac{m_0}{\sqrt{1 - \frac{V'^2}{c^2}}} = \frac{m_0}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$\text{Substitute for } M_0 = \frac{2m_0}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$\Rightarrow M' = \frac{2m_0}{\left(1 - \frac{v^2}{c^2}\right)}$$

Now

$$\begin{aligned} P'_{\text{before}} &= m'_1 v'_1 + m'_2 v'_2 \\ &= m_0(0) + m_0 \left(\frac{1+v^2/c^2}{1-v^2/c^2} \right) \left(\frac{-2v}{1+v^2/c^2} \right) \\ &= \frac{-2m_0v}{(1-v^2/c^2)} \end{aligned}$$

$$\begin{aligned} P'_{\text{after}} &= M' V' = \frac{2m_0}{(1-\frac{v^2}{c^2})} (-v) \\ &= \frac{-2m_0v}{(1-\frac{v^2}{c^2})} \end{aligned}$$

\Rightarrow

$$P'_{\text{before}} = P'_{\text{after}}$$

\Rightarrow

Momentum is conserved in S'

Summary:

$$M = \frac{m_0}{\sqrt{1 - \frac{v^2}{c^2}}} \quad \text{relativistic mass}$$

$$P = m_0 v = \frac{m_0 v}{\sqrt{1 - \frac{v^2}{c^2}}} \quad \text{relativistic momentum}$$

Calculate now ΔK

Assuming that the object starts from rest

\Rightarrow

$$\Delta K = K_f - K_i = K - 0 = K.$$

From work-energy theorem

$$\Delta K = K = W = \int F dx = \int \frac{dP}{dt} dx$$

\Rightarrow

$$K = \int \frac{dP}{dt} dx = \int v dp.$$

Integrating by parts $d(pv) = \underbrace{v dp}_{\text{---}} + p dv$

\Rightarrow

$$K = \int d(pv) - \int \overset{C_1 = v}{p} dv$$

$$= pv - \int_{C_1=0}^{C_2=v} p dv$$

$$= \frac{m_0 v}{\sqrt{1 - \frac{v^2}{c^2}}} v - \int_0^v \frac{m_0 v}{\sqrt{1 - \frac{v^2}{c^2}}} dv$$

$$= \frac{m_0 v^2}{\sqrt{1 - \frac{v^2}{c^2}}} - m_0 \int_0^v \frac{v dv}{\sqrt{1 - \frac{v^2}{c^2}}}$$

Change variables $v = c \cos \theta$

$$\Rightarrow dv = -c \sin \theta d\theta$$

$$\Rightarrow -m_0 \int_0^v \frac{v dv}{\sqrt{1 - \frac{v^2}{c^2}}} = +m_0 c^2 \int \frac{\sin \theta \cos \theta}{\sqrt{1 - \cos^2 \theta}} d\theta$$

$$-m_0 \int \frac{v dr}{\sqrt{1 - \frac{v^2}{c^2}}} = +m_0 c^2 \int \cos \theta d\theta$$

$$= +m_0 c^2 [\sin \theta]$$

$$= +m_0 c^2 [\sqrt{1 - \cos^2 \theta}]$$

$$= +m_0 c^2 \left[\sqrt{1 - \frac{v^2}{c^2}} \right]^0$$

$$= m_0 c^2 \sqrt{1 - \frac{v^2}{c^2}} - m_0 c^2$$

 \Rightarrow

$$K = \frac{m_0 v^2}{\sqrt{1 - \frac{v^2}{c^2}}} + m_0 c^2 \sqrt{1 - \frac{v^2}{c^2}} - m_0 c^2$$

$$= \frac{m_0 c^2}{\sqrt{1 - \frac{v^2}{c^2}}} \left[\frac{v^2}{c^2} + \left(1 - \frac{v^2}{c^2} \right) \right] - m_0 c^2$$

$$K = m c^2 - m_0 c^2$$

relativistic
Kinetic energy.

 \Rightarrow

$$K = m c^2 - E_0, \quad E_0 = \text{rest mass energy.}$$

or

$$E = K + E_0 = m c^2, \quad E = \text{total energy.}$$

$$E = m c^2$$

$$m = \frac{m_0}{\sqrt{1 - \frac{v^2}{c^2}}}$$

Go back to the completely inelastic collision problem and see if the total energy is conserved in the collision.

In S:

$$E_{m_1} = m_1 c^2 = \frac{m_0 c^2}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$E_{m_2} = m_2 c^2 = \frac{m_0 c^2}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$E_M = M_0 c^2 = \frac{2 m_0 c^2}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$E_{m_1} + E_{m_2} = E_M$$

Total energy is conserved in S.

In S':

$$E'_{m_1} = m_0 c^2$$

$$E'_{m_2} = m_0 c^2 \left(\frac{1 + \frac{v^2}{c^2}}{1 - \frac{v^2}{c^2}} \right)$$

$$E'_M = \frac{2 m_0 c^2}{(1 - v^2/c^2)}$$

$$\begin{aligned}
 \Rightarrow E_{m_1} + E_{m_2} &= m_0 c^2 + m_0 c^2 \left(\frac{1 + v^2/c^2}{1 - v^2/c^2} \right) \\
 &= m_0 c^2 \left[1 + \frac{\left(1 + \frac{v^2}{c^2}\right)}{\left(1 - \frac{v^2}{c^2}\right)} \right] \\
 &= m_0 c^2 \left[\frac{1 - \frac{v^2}{c^2} + 1 + \frac{v^2}{c^2}}{1 - \frac{v^2}{c^2}} \right] \\
 &= \frac{2 m_0 c^2}{\left(1 - \frac{v^2}{c^2}\right)}
 \end{aligned}$$

\Rightarrow

E_{total} is conserved in S

We can also derive an important relation between relativistic total energy and relativistic momentum which is

$$E^2 = p^2 c^2 + (m_0 c^2)^2$$

Proof:

$$p = \frac{m_0 v}{\sqrt{1 - \frac{v^2}{c^2}}} , \quad m = \frac{m_0}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$E = m c^2$$

\Rightarrow

$$\begin{aligned}
 p^2 c^2 + (m_0 c^2)^2 &= \frac{m_0^2 v^2 c^2}{\left(1 - \frac{v^2}{c^2}\right)} + m_0^2 c^4 \\
 &= \frac{m_0 v^2 c^2 + m_0^2 c^4 \left(1 - \frac{v^2}{c^2}\right)}{\left(1 - \frac{v^2}{c^2}\right)}
 \end{aligned}$$

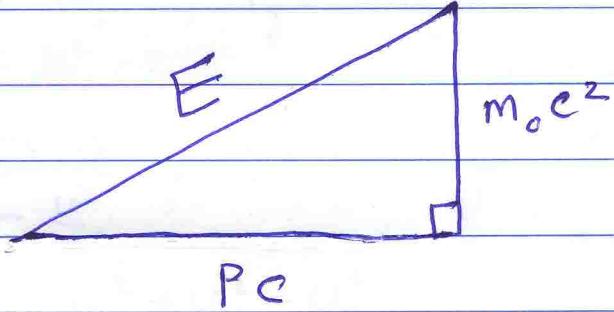
$$P^2 c^2 + (m_0 c^2)^2 = \frac{m_0^2 v^2 c^2 + m_0^2 c^4 - m_0^2 v^2 c^2}{(1 - \frac{v^2}{c^2})}$$

$$= \left(\frac{m_0 c^2}{\sqrt{1 - \frac{v^2}{c^2}}} \right)^2$$

$$= (mc^2)^2$$

$$= E^2$$

\Rightarrow $E^2 = P^2 c^2 + (m_0 c^2)^2$



Next we consider the following

(P_x, P_y, P_z, E) as the components of momentum and total energy of a particle in S frame

This is called (momentum-Energy 4-vector)

Remember that we learned about space-time 4-vector (x, y, z, t)

Now in S', what is the momentum-energy 4-vector?

Solution: $S' \rightarrow v$ along +ve x-axis.

In S:

$$P_x = m u_x = \frac{m_0 u_x}{\sqrt{1 - \frac{u^2}{c^2}}}, \quad P_y = m u_y = \frac{m_0 u_y}{\sqrt{1 - \frac{u^2}{c^2}}}$$

$$P_z = m u_z = \frac{m_0 u_z}{\sqrt{1 - \frac{u^2}{c^2}}}, \quad E = mc^2 = \frac{m_0 c^2}{\sqrt{1 - \frac{u^2}{c^2}}}$$

Of course $u^2 = u_x^2 + u_y^2 + u_z^2$

In S':

$$P'_x = m' u'_x = \frac{m_0 u_x}{\sqrt{1 - \frac{u'^2}{c^2}}}, \quad P'_y = m' u'_y = \frac{m_0 u_y}{\sqrt{1 - \frac{u'^2}{c^2}}}$$

$$P'_z = m' u'_z = \frac{m_0 u'_z}{\sqrt{1 - \frac{u'^2}{c^2}}}, \quad E' = m' c^2 = \frac{m_0 c^2}{\sqrt{1 - \frac{u'^2}{c^2}}}$$

Of course $u'^2 = u'_x^2 + u'_y^2 + u'_z^2$



$$17/1 \quad u^2 = \frac{(u_x - v)^2}{(1 - \frac{u_x v}{c^2})^2} + \frac{(u_y \sqrt{1 - v^2/c^2})^2}{(1 - \frac{u_x v}{c^2})^2} + \frac{(u_z \sqrt{1 - v^2/c^2})^2}{(1 - \frac{u_x v}{c^2})^2} \quad 13$$

$$= \frac{(u_x - v)^2}{(1 - \frac{u_x v}{c^2})^2} + u_y^2 (1 - \frac{v^2}{c^2}) + u_z^2 (1 - \frac{v^2}{c^2})$$

$$(1 - \frac{u^2}{c^2}) = 1 - \frac{(\frac{u_x}{c} - \frac{v}{c})^2 + (\frac{u^2}{c^2} - \frac{u_x^2}{c^2})(1 - \frac{v^2}{c^2})}{(1 - \frac{u_x v}{c^2})^2}$$

$$= \frac{(1 - \frac{u_x v}{c^2})^2 - (\frac{u_x}{c} - \frac{v}{c})^2 - (\frac{u^2}{c^2} - \frac{u_x^2}{c^2})(1 - \frac{v^2}{c^2})}{(1 - \frac{u_x v}{c^2})^2}$$

$$= \frac{1 - 2 \frac{u_x v}{c^2} + \frac{u_x^2 v^2}{c^4} - \frac{u_x^2}{c^2} - \frac{v^2}{c^2} + 2 \frac{u_x v}{c^2}}{(1 - \frac{u_x v}{c^2})^2}$$

$$- \frac{u^2}{c^2} + \frac{u_x^2}{c^2} + \frac{u^2 v^2}{c^4} - \frac{u_x^2 v^2}{c^4}$$

$$= \frac{(1 - \frac{v^2}{c^2}) - \frac{u^2}{c^2}(1 - \frac{v^2}{c^2})}{(1 - \frac{u_x v}{c^2})^2}$$

$$= \frac{(1 - \frac{v^2}{c^2})(1 - \frac{u^2}{c^2})}{(1 - \frac{u_x v}{c^2})^2}$$

$$\Rightarrow \sqrt{1 - \frac{u^2}{c^2}} = \sqrt{\frac{(1 - \frac{v^2}{c^2})(1 - \frac{u^2}{c^2})}{(1 - \frac{u_x v}{c^2})^2}}$$

$$= \frac{\sqrt{1 - v^2/c^2} \sqrt{1 - u^2/c^2}}{(1 - \frac{u_x v}{c^2})}$$

$$\sqrt{1 - \frac{u'^2}{c^2}} = \frac{\sqrt{1 - u^2/c^2}}{\gamma \left(1 - \frac{u_x v}{c^2}\right)}$$

Similarly we can show that

$$\sqrt{1 - \frac{u^2}{c^2}} = \frac{\sqrt{1 - u'^2/c^2}}{\gamma \left(1 + \frac{u_x v}{c^2}\right)}$$

Go back to our problem (momentum-energy 4-vector)

In S':

$$\vec{P}'_x = \frac{m_0 u'_x}{\sqrt{1 - \frac{u'^2}{c^2}}} , \quad u'_x = \frac{u_x - v}{\sqrt{1 - \frac{u_x v}{c^2}}}$$

and

$$\frac{1}{\sqrt{1 - \frac{u'^2}{c^2}}} = \frac{\left(1 - \frac{u_x v}{c^2}\right)}{\sqrt{1 - \frac{v^2}{c^2}} \sqrt{1 - \frac{u^2}{c^2}}}$$

$$\Rightarrow \vec{P}'_x = m_0 \frac{(u_x - v)}{\left(1 - \frac{u_x v}{c^2}\right)} \frac{\left(1 - \frac{u_x v}{c^2}\right)}{\sqrt{1 - \frac{v^2}{c^2}} \sqrt{1 - \frac{u^2}{c^2}}}$$

$$= \frac{m_0}{\sqrt{1 - \frac{u^2}{c^2}}} \frac{(u_x - v)}{\sqrt{1 - \frac{v^2}{c^2}}} \quad \gamma = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$= \gamma (m u_x - m v)$$

$$= \gamma \left(P_x - \frac{E_0 v}{c^2} \right).$$

Similarly, we can show that

$$\vec{P}'_y = P_y, \quad \vec{P}'_z = P_z$$

Also

$$E' = m'c^2 = \frac{m_0 c^2}{\sqrt{1 - \frac{u'^2}{c^2}}}$$

$$= \frac{m_0 c^2}{\sqrt{1 - \frac{u^2}{c^2}}} \cdot \frac{\left(1 - \frac{u_x v}{c^2}\right)}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$= \gamma m c^2 \left(1 - \frac{u_x v}{c^2}\right)$$

$$(E' = \gamma(E - P_x v))$$

The inverse relation is

$$E = \gamma(E' + P_x v)$$

Note that (P_x, P_y, P_z) transform like (x, y, z) and E transforms like t .

Also

$$P'^2 - \frac{E'^2}{c^2} = P_x'^2 + P_y'^2 + P_z'^2 - \frac{E'^2}{c^2}$$

$$= \gamma^2 \left(P_x - \frac{E \gamma}{c^2} \right)^2 + P_y^2 + P_z^2 - \frac{\gamma^2}{c^2} (E - P_x \gamma)^2$$

$$= \gamma^2 \left(P_x^2 - 2 \frac{P_x E \gamma}{c^2} + \frac{E^2 \gamma^2}{c^4} \right) + P_y^2 + P_z^2$$

$$- \gamma^2 \left(\frac{E^2}{c^2} - 2 \frac{P_x E \gamma}{c^2} + \frac{P_x^2 \gamma^2}{c^2} \right)$$

$$= \cancel{\gamma^2 P_x^2} - \cancel{\gamma^2 \left(2 \frac{P_x E \gamma}{c^2} \right)} + \cancel{\gamma^2 \left(\frac{E^2 \gamma^2}{c^4} \right)} + P_y^2 + P_z^2$$

$$- \cancel{\gamma^2 \frac{E^2}{c^2}} + \cancel{\gamma^2 \left(2 \frac{P_x E \gamma}{c^2} \right)} - \cancel{\gamma^2 \frac{P_x^2 \gamma^2}{c^2}}$$

$$= \gamma^2 P_x^2 \left(1 - \frac{\gamma^2}{c^2} \right) - \gamma^2 \frac{E^2}{c^2} \left(1 - \frac{\gamma^2}{c^2} \right)$$

$$+ P_y^2 + P_z^2$$

$$\text{But } \gamma^2 \left(1 - \frac{\gamma^2}{c^2} \right) = 1.$$

⇒

$$P'^2 - \frac{E'^2}{c^2} = P_x^2 + P_y^2 + P_z^2 - \frac{E^2}{c^2}$$

or

$$\boxed{P'^2 - \frac{E'^2}{c^2} = P^2 - \frac{E^2}{c^2}}$$

Transformation of forces in different inertial frames

In frame S:

$$\vec{F} = F_x \hat{i} + F_y \hat{j} + F_z \hat{k}$$

In frame S':

$$\vec{F}' = \frac{d\vec{P}'}{dt'} = \frac{\left(\frac{d\vec{P}'}{dt}\right)}{\left(\frac{dt'}{dt}\right)}$$

$$F'_x = \frac{\left(\frac{dP'_x}{dt}\right)}{\left(\frac{dt'}{dt}\right)}$$

$$P'_x = \gamma \left(P_x - \frac{E u}{c^2} \right)$$

$$t' = \gamma \left(t - \frac{u x}{c^2} \right)$$

$$\frac{dP'_x}{dt} = \gamma \left(\frac{dP_x}{dt} - \frac{u}{c^2} \frac{dE}{dt} \right)$$

$$\frac{dt'}{dt} = \gamma \left(1 - \frac{u}{c^2} \frac{dx}{dt} \right)$$

Substitute \Rightarrow

$$F'_x = \gamma \left(\frac{dP_x}{dt} - \frac{u}{c^2} \frac{dE}{dt} \right)$$

$$\gamma \left(1 - \frac{u}{c^2} u_x \right)$$

$$F'_x = F_x - \frac{u}{c^2} \frac{dE}{dt} \quad (1)$$

$$\left(1 - \frac{u}{c^2} u_x \right)$$

$$\text{Since } E^2 = P^2 c^2 + m_0^2 c^4 = c^2 (\vec{P} \cdot \vec{P}) + m_0^2 c^4$$

$$\Rightarrow 2E \frac{dE}{dt} = c^2 \vec{P} \cdot \frac{d\vec{P}}{dt} + c^2 \frac{d\vec{P}}{dt} \cdot \vec{P}$$

$$\cancel{2E \frac{dE}{dt}} = \cancel{2c^2 \vec{P} \cdot \frac{d\vec{P}}{dt}}, \quad E = mc^2$$

$$\frac{dE}{dt} = \frac{\vec{P}}{m} \cdot \frac{d\vec{P}}{dt} = \vec{F} \cdot \vec{u}$$

Substitute in (v) \Rightarrow

$$\vec{F}'_x = \underline{\underline{F}_x - \frac{v}{c^2} (\vec{F} \cdot \vec{u})} \\ (1 - \frac{v}{c^2} u_x)$$

Similarly \Rightarrow

$$\vec{F}'_y = \underline{\underline{F}_y} \\ \gamma (1 - \frac{v}{c^2} u_x)$$

$$\vec{F}'_z = \underline{\underline{F}_z} \\ \gamma (1 - \frac{v}{c^2} u_x)$$

The inverse relations are: $v \rightarrow -v$

$$\vec{F}_x = \underline{\underline{\vec{F}'_x + \frac{v}{c^2} (\vec{F}' \cdot \vec{u}')}} \\ (1 + \frac{v}{c^2} u'_x)$$

$$\vec{F}_y = \underline{\underline{\vec{F}'_y}} \\ \gamma (1 + \frac{v}{c^2} u'_x)$$

$$\vec{F}_z = \underline{\underline{\vec{F}'_z}} \\ \gamma (1 + \frac{v}{c^2} u'_x)$$

If the particle in frame S is at rest $\Rightarrow \vec{u} = 0$

$$\Rightarrow \vec{F}'_x = \vec{F}_x$$

$$\vec{F}'_y = \underline{\underline{\frac{\vec{F}_y}{\gamma}}} \\ \gamma$$

$$\vec{F}'_z = \underline{\underline{\frac{\vec{F}_z}{\gamma}}} \\ \gamma$$

Clarify two basic concepts in Physics

- * Concept of invariance
- * Concept of conservation.

- * A quantity is invariant if it is the same in different inertial frames.

Example: Rest mass of a particle.

- * A quantity is conserved if it remains unchanged in time, i.e. it remains the same after some interaction process.

Example: Linear momentum, Total energy

- * Invariant quantities, such as rest mass, are not necessarily conserved.

- * Conserved quantities, such as total energy and momentum, are not invariant since they change from one inertial frame to another.

Example: The quantity $E^2 - p^2 c^2$

$$\text{Since } E^2 = p^2 c^2 + E_0^2 \quad E_0 = m_0 c^2$$

$$\Rightarrow E^2 - p^2 c^2 = E_0^2 = E'^2 - p'^2 c^2$$

i.e. the quantity $E^2 - p^2 c^2$ is an invariant quantity.

Photons:

- * Light is a wave that carries energy and exerts a pressure when it is incident on a surface or object.
- * Photons are considered as massless particles.
- * From $E^2 = p^2 c^2 + m_0^2 c^4 \Rightarrow$
For photons
- * Energy carried by photons $E = h\nu$, where h is Planck's constant, ν = frequency.

$$P = \frac{h\nu}{c}$$

But $c = \nu\lambda \Rightarrow \frac{c}{\nu} = \lambda$.

$$P = \frac{h}{\lambda}$$