## Steady-State Temperature in a Sphere

## Laplace's equation in spherical coordinates

## Problem 1:

Find the steady-state temperature inside a sphere of radius r = 1when the surface of the upper half is held at 100° and the surface of the lower half at 0°.

## Solution:

- a) Since there is no source of heat is available inside the sphere, the temperature *u* satisfies Laplace's equation.
- b) The symmetry of the problem implies the use of spherical coordinates.

$$\nabla^2 u = \frac{1}{h_1 h_2 h_3} \left[ \frac{\partial}{\partial x_1} \left( \frac{h_2 h_3}{h_1} \frac{\partial u}{\partial x_1} \right) + \frac{\partial}{\partial x_2} \left( \frac{h_1 h_3}{h_2} \frac{\partial u}{\partial x_2} \right) + \frac{\partial}{\partial x_3} \left( \frac{h_1 h_2}{h_3} \frac{\partial u}{\partial x_3} \right) \right]$$

In spherical coordinates:

$$h_1 = 1 \qquad h_2 = r \quad h_3 = r \sin \theta$$

$$\partial x_1 = \partial r \qquad \partial x_2 = \partial \theta \qquad \partial x_3 = \partial \phi$$
$$\nabla^2 u = \frac{1}{r^2 \sin \theta} \left[ \frac{\partial}{\partial r} \left( r^2 \sin \theta \frac{\partial u}{\partial r} \right) + \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial u}{\partial \theta} \right) + \frac{\partial}{\partial \phi} \left( \frac{1}{\sin \theta} \frac{\partial u}{\partial \phi} \right) \right]$$

But

$$\nabla^2 u = 0$$

$$\therefore \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial u}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial u}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial u^2}{\partial^2 \phi} = 0$$

Try a solution of type  $u(r, \theta, \phi) = R(r)\Theta(\theta)\Phi(\phi)$ 

Substitute this solution into the PDF and multiply both sides by

$$\frac{r^{2}\sin^{2}\theta}{R\Theta\Phi} \text{ to get:}$$

$$\frac{\sin^{2}\theta}{R}\frac{d}{dr}(r^{2}\frac{dR}{dr}) + \frac{\sin\theta}{\Theta}\frac{d}{d\theta}(\sin\theta\frac{d\Theta}{d\theta}) + \frac{1}{\Phi}\frac{d^{2}\Phi}{d\phi^{2}} = 0$$

Put  $\frac{1}{\Phi} \frac{d^2 \Phi}{d\phi^2} = -m^2$  (The solution to this DE is  $\Phi = C \sin \phi + D \cos \phi$ )

After substituting the latter DF into the PDF and dividing by  $sin^2\theta$ , the PDF becomes:

$$\frac{1}{R}\frac{d}{dr}(r^2\frac{dR}{dr}) + \frac{1}{\Theta\sin\theta}\frac{d}{d\theta}(\sin\theta\frac{d\Theta}{d\theta}) - \frac{m}{\sin^2\theta} = 0$$

Now this DF is separable. The radial part of this equation is set equal to a constant k.

$$\frac{1}{R}\frac{d}{dr}(r^2\frac{dR}{dr}) = k$$

[Note: It is more suitable to write *k* as the product of two successive integers (*i.e.*  $k = \ell(\ell + 1)$ )].

$$\frac{d}{dr}(r^2\frac{dR}{dr}) = \ell(\ell+1)R$$

The last differential equation has the form:

$$r^{2} \frac{d^{2}R}{dr^{2}} + 2r \frac{dR}{dr} - \ell(\ell+1)R = 0$$

Assume a solution:  $R = r^n$  and substitute it into the DE to get:

$$n(n-1) r^{n} + 2nr^{n} - \ell(\ell+1) r^{n} = 0$$

Equating the coefficients of  $r^n$  in this equation

$$n^2 + n - \ell(\ell + 1) = 0$$

Thus *n* has two roots:

$$n = \ell$$
 and  $n = -(\ell + 1)$ 

The general solution to the radial equation is a linear combination

of two solutions, *i.e.*  $R = Ar^{\ell} + Br^{-\ell-1}$ 

Remainder of DF:

$$\frac{1}{\sin\theta} \frac{d}{d\theta} (\sin\theta \frac{d\Theta}{d\theta}) + \left[\ell(\ell+1) - \frac{m}{\sin^2\theta}\right] \Theta = 0$$

[This is the associated Legendre DE which hives the common solution of associated Legendre polynomial, i.e.  $\Theta = P_{\ell}^{m}(\cos\theta)$  (see problem 10.2 in chapter 12).

Thus the general solution  $(u = R\Theta\Phi)$  becomes:

$$u = (Ar^{\ell} + Br^{-\ell-1})(C\sin m\phi + D\cos m\phi)P_{\ell}^{m}(\cos\theta)$$

Notes:

- 1) Since we are interested to find the temperature inside the sphere, we have to consider B = 0, because  $r^{-\ell-1}$  goes to infinity at the origin (r = 0).
- 2) The problem has azimuthal symmetry i.e. u is independent of φ (as φ changes u is constant). This implies that u = D cos mφ with m = 0 and cos mφ.

⇒ u = D [This can be justified for the given B.C where the top of the sphere is at 100° and the bottom of the sphere at 0°]. For  $m \neq 0$ 

 $u_{\ell} = A'r^{\ell}\cos m\phi P_{\ell}^{m}(\cos\theta)$ 

[Note: The spherical harmonic  $Y_{\ell}^{m}(\theta, \phi)$  is related to the associated Legendre polynomial can be expressed

as: 
$$Y_{\ell}^{m}(\theta,\phi)\sqrt{\frac{2\pi}{2\ell+}\frac{(\ell-m)!}{(\ell+m)!}} = P_{\ell}^{m}(\cos\theta)\cos m\phi$$
].

Here, in this problem we have m = 0 and the solution is reduced

to: 
$$u_{\ell} = A' r^{\ell} P_{\ell}^{m}(\cos \theta)$$
  
 $\Rightarrow u = \sum_{\ell} u_{\ell}$ .  
And  $u = \sum_{\ell=0}^{\infty} A'_{\ell} r^{\ell} P_{\ell}(\cos \theta)$ 

The coefficients  $A'_{\ell}$  can be determined using the given temperatures when r = 1.

$$u(r=1,\theta) = \begin{cases} 100 & 0 < \theta < \frac{\pi}{2} \\ 0 & \frac{\pi}{2} < \theta < \pi \end{cases}$$

$$u(r=1, heta)=\sum_{\ell=0}^{\infty}A_{\ell}'P_{\ell}(x)$$
 , (where  $x=\cos heta$ )

Also we have  $u(r=1,\theta) = 100f(x)$ , where

$$f(x) = \begin{cases} 0 & -1 < x < 0 \\ 1 & 0 < x < 1 \\ \therefore \sum_{\ell=0}^{\infty} A_{\ell}' P_{\ell}(x) = 100 f(x) \\ A_{\ell}' = \frac{2\ell + 1}{2} \int_{-1}^{1} u(x) P_{\ell}(x) dx \\ \Rightarrow A_{\ell}' = \frac{2\ell + 1}{2} [100 \int_{0}^{1} P_{\ell}(x) dx] \\ \text{Thus } A_{0} = \frac{100}{2} \int_{0}^{1} dx = \frac{100}{2}; A_{1} = \frac{300}{2} \int_{0}^{1} x dx = \frac{300}{4} \text{ and } A_{3} = -\frac{700}{16} \dots \text{ect.} \\ \therefore u(1, \theta) = 100 [\frac{1}{2} P_{0}(\cos \theta) + \frac{3}{4} P_{1}(\cos \theta) - \frac{7}{16} P_{3}(\cos \theta) + \bullet \bullet \bullet]. \end{cases}$$
(Solve problem 7.13)