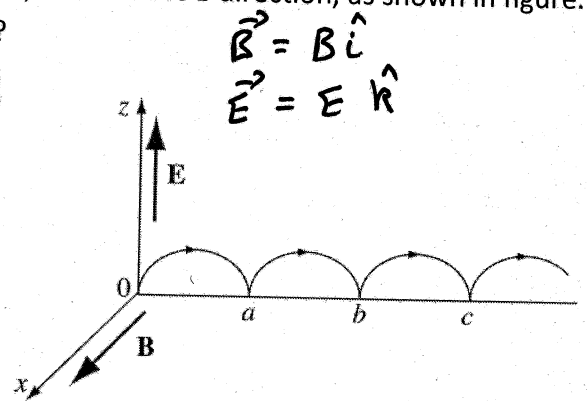


Example 5.2. Cycloid Motion. A more exotic trajectory occurs if we include a uniform electric field, at right angles to the magnetic one. Suppose, for instance, that \mathbf{B} points in the x-direction, and \mathbf{E} in the z-direction, as shown in figure. A positive charge (q) is released from the origin; what path will it follow?

Let's think it through qualitatively, first. Initially, the particle is at rest, so the magnetic force is zero, and the electric field accelerates the charge in the z-direction. As it picks up speed, a magnetic force develops which pulls the charge around to the right. The faster it goes, the stronger F_{mag} becomes; eventually, it curves the particle back around towards the y-axis. At this point the charge is moving against the electrical force, so it begins to slow down; the magnetic force then decreases, and the electrical force takes over, bringing the particle to rest at point a , in figure. There the entire process commences anew, carrying the particle over to point b , and so on.



$$\vec{B} = B \hat{i}$$

$$\vec{E} = E \hat{k}$$

⇒ No motion along x-direction
⇒ $U_x = 0$

$$\vec{F} = q\vec{E} + q\vec{v} \times \vec{B}$$

$$= qE\hat{k} + qBv_z\hat{j} - qBv_y\hat{k}$$

$$= qBv_z\hat{j} + q(E - Bv_y)\hat{k}$$

as seen no force along the x-direction

$$\vec{v} = (0, v_y, v_z) \Rightarrow m \frac{d\vec{v}}{dt} = \vec{F}$$

$$\Rightarrow m \frac{dv_y}{dt} \hat{j} + m \frac{dv_z}{dt} \hat{k} = qBv_z\hat{j} + q(E - Bv_y)\hat{k}$$

$$\Rightarrow m \frac{dv_y}{dt} = qBv_z \Rightarrow \frac{dv_y}{dt} = \frac{qB}{m} v_z = \omega v_z \quad \text{--- (1)}$$

$$\text{and } m \frac{dv_z}{dt} = q(E - Bv_y) \Rightarrow \frac{dv_z}{dt} = \frac{qE}{m} - \frac{qB}{m} v_y \quad \text{--- (2)}$$

(1) and (2) are two coupled differential equations, they are easily solved by differentiating the first and using the second to eliminate v_z ;

$$\dot{U}_y = \omega U_z \quad \text{--- (1)}$$

$$\text{where } \omega = \frac{qB}{m}$$

$$\dot{U}_z = \frac{qE}{m} - \omega U_y \quad \text{--- (2)}$$

so differentiating (1), $\ddot{U}_y = \omega \dot{U}_z = \omega \left(\frac{qE}{m} - \omega U_y \right)$

$$\Rightarrow \boxed{\ddot{U}_y + \omega^2 U_y = \frac{qE\omega}{m}} \quad \text{--- (3) non-homogeneous 2nd order linear diff. eqⁿ}$$

similarly differentiate (2) and use (1) to eliminate \dot{U}_y

$$\ddot{U}_z = -\omega \dot{U}_y = -\omega(\omega U_z) = -\omega^2 U_z$$

$$\Rightarrow \boxed{\ddot{U}_z + \omega^2 U_z = 0} \quad \text{--- (4) homogeneous 2nd order diff eqⁿ}$$

The solution of (3) is composed of two parts: complementary solution and particular solution

- comp. solution $\ddot{U}_y + \omega^2 U_y = 0 \Rightarrow m^2 + \omega^2 = 0$
auxiliary eqⁿ

$$U_y^c(b) = c_1 e^{m_1 b} + c_2 e^{m_2 b}$$

$$= c_1 e^{i\omega b} + c_2 e^{-i\omega b}$$

$$= c_1 \cos \omega b + i c_1 \sin \omega b + c_2 \cos \omega b - c_2 i \sin \omega b$$

$$= (c_1 + c_2) \cos \omega b + (c_1 - c_2) i \sin \omega b$$

$$U_y^c(b) = A_1 \cos \omega b + A_2 \sin \omega b$$

$$m^2 = -\omega^2$$

$$m = \pm \sqrt{-\omega^2} = \pm \sqrt{-1} \omega$$

$$= \pm i\omega$$

- particular solution: since the term in the R.H.S is constant, then $U_y^p(b) = \text{constant} = D$; in general the particular solution takes the same form of what exists on the R.H.S, for example $ay'' + by' + cy = kx^2$, then $y^p(x) = ax^2 + bx + c$, a polynomial of the same degree

$\therefore U_y(t) = D$; substitute this into (3), we get

$$\ddot{U}_y + \omega^2 U_y = \frac{qEw}{m} \Rightarrow 0 + \omega^2 D = \frac{qEw}{m} \Rightarrow D = \frac{qEw}{m\omega^2} = \frac{qE}{m\omega}$$

$$\Rightarrow D = \frac{qE}{m\omega} = \frac{E}{B}$$

$$\therefore U_y(t) = U_y^c(t) + U_y^p(t)$$

$$U_y(t) = A_1 \cos \omega t + A_2 \sin \omega t + \frac{E}{B} \quad \dots (5)$$

the solution of equation (4) is

$$U_z(t) = B_1 \cos \omega t + B_2 \sin \omega t \quad \dots (6)$$

now differentiate (5) and equate the result with (4)

$$\dot{U}_y = -A_1 \omega \sin \omega t + A_2 \omega \cos \omega t = \omega U_z$$

$$\Rightarrow -A_1 \omega \sin \omega t + A_2 \omega \cos \omega t = \omega B_1 \cos \omega t + \omega B_2 \sin \omega t$$

$$\Rightarrow A_2 \omega = \omega B_1 \Rightarrow A_2 = B_1$$

$$\text{and } -A_1 \omega = \omega B_2 \Rightarrow A_1 = -B_2$$

so eq^s (5) and (6) become

$$\Rightarrow U_y = A_1 \cos \omega t + A_2 \sin \omega t + \frac{E}{B} \quad \dots (7)$$

$$U_z = A_2 \cos \omega t - A_1 \sin \omega t \quad \dots (8)$$

integrating (7) and (8), we get

$$y(t) = \frac{A_1}{\omega} \sin \omega t - \frac{A_2}{\omega} \cos \omega t + \frac{E}{B} t + A_3 \quad \dots (9)$$

$$z(t) = \frac{A_2}{\omega} \sin \omega t + \frac{A_1}{\omega} \cos \omega t + A_4 \quad \dots (10)$$

now A_1, A_2, A_3, A_4 are found from Boundary conditions

since the particle starts from rest at the origin
 $\Rightarrow v_y(b=0) = v_z(b=0) = 0$ and $y(b=0) = z(b=0) = 0$
 so from v_y and v_z equations (7) and (8), we get

$$0 = A_1 + \frac{E}{B} \Rightarrow \boxed{A_1 = -\frac{E}{B}} \text{ and } \boxed{0 = A_2}$$

$$\therefore y(b) = -\frac{E}{B\omega} \sin \omega b + \frac{E}{B} t + A_3 \quad \text{--- (11)}$$

$$z(b) = \frac{E}{B\omega} \cos \omega b + A_4 \quad \text{--- (12)}$$

Now from the second B.Cs ($y(b=0) = z(b=0) = 0$)

we get

$$0 = A_3 \text{ and } 0 = -\frac{E}{B\omega} + A_4 \Rightarrow A_4 = +\frac{E}{B\omega}$$

$$\text{so } y(b) = -\frac{E}{B\omega} \sin \omega b + \frac{E}{B} t = \frac{E}{B\omega} (\omega b - \sin \omega b) \quad \text{--- (13)}$$

$$\text{and } z(b) = -\frac{E}{B\omega} \cos \omega b + \frac{E}{B\omega} = \frac{E}{B\omega} (1 - \cos \omega b) \quad \text{--- (14)}$$

define $\frac{E}{B\omega} = R$, radius of the circle, then

to find the combined equation of the path, take

$$y^2 + z^2 = R^2 (\omega b - \sin \omega b)^2 + R^2 (1 - \cos \omega b)^2$$

$$= R^2 [\omega^2 b^2 + \sin^2 \omega b - 2\omega b \sin \omega b]$$

$$+ R^2 [1 + \cos^2 \omega b - 2 \cos \omega b]$$

$$= R^2 \omega^2 b^2 + R^2 \sin^2 \omega b - 2\omega b R^2 \sin \omega b$$

$$+ R^2 + R^2 \cos^2 \omega b - 2 R^2 \cos \omega b$$

using $\sin^2 \omega b + \cos^2 \omega b = 1$, we get

$$y^2 + z^2 = R^2 \omega^2 b^2 + 2R^2 - 2\omega b R^2 \sin \omega b - 2R^2 \cos \omega b \quad \dots (15)$$

but from (13) and (14), we have

$$\left. \begin{aligned} \frac{y}{R} = \omega b - \sin \omega b &\Rightarrow -\sin \omega b = \frac{y}{R} - \omega b \\ \text{and } \frac{z}{R} = 1 - \cos \omega b &\Rightarrow -\cos \omega b = \frac{z}{R} - 1 \end{aligned} \right\} \text{ substitute in (15)}$$

$$\begin{aligned} \Rightarrow y^2 + z^2 &= R^2 \omega^2 b^2 + 2R^2 + 2\omega b R^2 \left(\frac{y}{R} - \omega b \right) + 2R^2 \left(\frac{z}{R} - 1 \right) \\ &= R^2 \omega^2 b^2 + \cancel{2R^2} + 2\omega b R y - 2\omega^2 b^2 R^2 + 2Rz - \cancel{2R^2} \\ &= R^2 \omega^2 b^2 + 2\omega b R y - 2\omega^2 b^2 R^2 + 2Rz \\ &= -R^2 \omega^2 b^2 + 2\omega b R y + 2Rz \end{aligned}$$

$$\Rightarrow \underbrace{y^2 + R^2 \omega^2 b^2 - 2\omega b R y}_{(y - R\omega b)^2} + z^2 - 2Rz = 0$$

$$\Rightarrow (y - R\omega b)^2 + z^2 - 2Rz = 0,$$

completing the square of $z^2 - 2Rz$ by adding $(R^2 - R^2)$, we get

$$(y - R\omega b)^2 + z^2 - 2Rz + R^2 = R^2$$

$$(y - R\omega b)^2 + (z - R)^2 = R^2$$

equation of a circle of radius R , whose center $(0, R\omega b, R)$ travels in the y -direction with a constant speed of $u = \omega R = \omega \frac{E}{B\omega} = \frac{E}{B}$. The curve of the motion is called cycloid and shown in the figure.